

ManyVal '08 - Applications of Topological Dualities to Measure Theory in Algebraic Many-Valued Logic, May 19–21, 2008, University of Milan, Milan, Italy

Coauthors: Zdena Riečanová, Department of Mathematics, Faculty of Electrical Engineering and Information Technology, Slovak University of Technology Ilkovičova 3, SK-812 19 Bratislava, Slovak Republic, zdena.riecanova@gmail.com

## LATTICE EFFECT ALGEBRAS POSSESING TWO-VALUED STATES

JAN PASEKA

Common generalizations of MV-algebras [2] and orthomodular lattices are lattice effect algebras [3]. An *effect algebra*  $(E; \oplus, 0, 1)$  is a set  $E$  with two special elements  $0, 1$  and a partial binary operation  $\oplus$  which is commutative and associative at which these equalities hold if one of their sides exists. Moreover, to every element  $a \in E$  there exists a unique element  $a' \in E$  with  $a \oplus a' = 1$  and if  $a \oplus 1$  exists then  $a = 0$ . In every effect algebra we can define a partial order by  $a \leq b$  iff there exists  $c \in E$  with  $a \oplus c = b$  (we set  $c = b \ominus a$ ). If  $(E; \leq)$  is a lattice (a complete lattice) then  $(E; \oplus, 0, 1)$  is called a *lattice effect algebra* (a *complete lattice effect algebra*).

Generalized effect algebras as posets are unbounded versions of effect algebras. In this case instead of the axiom on the existence of  $a'$  with  $a \oplus a' = 1$  for all  $a \in E$  we have cancellation law, i.e.,  $a \oplus b = a \oplus c$  implies  $b = c$  and, moreover,  $a \oplus b = 0$  implies  $a = b = 0$ .

A well known fact is that every generalized effect algebra  $P$  can be uniquely extended onto effect algebra  $E$  (called an *effect algebraic extension of  $P$* ) in which  $P$  is an order ideal in  $E$  and  $P^* = E \setminus P$  is a dual poset to  $P$ . We write  $E = P \dot{\cup} P^*$  (a disjoint union) [4]. On the other hand not every (lattice) effect algebra  $E$  becomes this way. We can prove

**Theorem 1.** [6] *Let  $(E; \oplus, 0, 1)$  be an effect algebra. The following conditions are equivalent:*

- (i) *There exists a two-valued state  $\omega$  on  $E$ .*
- (ii) *There exists a sub-generalized effect algebra  $P_\omega$  of  $E$  such that  $E = P_\omega \dot{\cup} P_\omega^*$ , where  $P_\omega^* = \{1 \ominus x \mid x \in P_\omega\}$ ,  $P_\omega \cap P_\omega^* = \emptyset$  and  $P_\omega$  is an order ideal in  $E$ .*

*In this case  $P_\omega = \omega^{-1}(\{0\})$ .*

For Archimedean atomic lattice effect algebra  $E$  we can show a sufficient condition (F) for the existence of a two-valued state  $\omega$  on  $E$ :

(F) There exists a finite set  $F = \{p_\kappa \mid \kappa \in H\}$  of pairwise noncompatible atoms of  $E$  such that for every atomic block  $M$  of  $E$  there exists  $\kappa_M \in H$  such that  $p_{\kappa_M} \in C(M)$  and  $\omega(p_{\kappa_M}) = 1$ . [6]

Moreover, we can prove that this condition (F) is a necessary and sufficient condition for the existence of a two-valued state  $\omega$  on: (a) every (o)-continuous Archimedean atomic lattice effect algebra, (b) every block-finite Archimedean atomic lattice effect algebra. Here a *block* of a lattice effect algebra is a maximal sub-lattice effect algebra being an MV-algebra. A lattice effect algebra is called *block-finite* if it has only finitely many blocks.

Using this facts we can prove

**Theorem 2.** [7] *Every Archimedean atomic lattice effect algebra with at most five blocks possess a state.*

Note that in present time the known example of finite lattice effect algebra admitting no states has nineteen blocks.

Finally, since the existence of a two-valued state  $\omega$  on an effect algebra  $E$  is equivalent to the fact that  $E$  is an effect algebraic extension of a sub-generalized effect algebra  $P_\omega = \{x \in E \mid \omega(x) = 0\}$ , the following question arises:

If  $F_1$  and  $F_2$  are two sets of pairwise noncompatible atoms of an Archimedean atomic lattice effect algebra  $E$  and  $\omega_1, \omega_2$  are two-valued states on  $E$  at which  $P_{F_1} = \omega_1^{-1}(\{0\})$  and  $P_{F_2} = \omega_2^{-1}(\{0\})$  are sub-generalized effect algebras of  $E$  with  $E = P_{F_1} \dot{\cup} P_{F_1}^* = P_{F_2} \dot{\cup} P_{F_2}^*$  whether (or at which conditions)  $P_{F_1}$  and  $P_{F_2}$  are isomorphic generalized effect algebras.

We show that two non-isomorphic generalized effect algebras  $P_{F_1}$  and  $P_{F_2}$  may have a common (or isomorphic) effect algebraic extension.

In spite of the fact that an atomic lattice effect algebra  $E$  may have nonatomic block [1] the following statement can be proved:

**Theorem 3.** [5] *Let  $E$  be an Archimedean atomic lattice effect algebra and  $F_1, F_2$  are two sets of pairwise noncompatible atoms satisfying the condition (F). Let there exists a bijection  $\psi : F_1 \rightarrow F_2$  such that, for every atomic block  $M$  of  $E$  we have  $a \in M \cap F_1 \iff \psi(a) \in M \cap F_2$ . Then  $P_{F_1} \cong P_{F_2}$ .*

This very simple condition is not necessary for the isomorphism of  $P_{F_1}$  and  $P_{F_2}$ . The necessary and sufficient condition for  $P_{F_1}$  and  $P_{F_2}$  is the following:

**Theorem 4.** [5] *Let  $E_1, E_2$  be Archimedean atomic lattice effect algebras,  $F_1 \subseteq E_1, F_2 \subseteq E_2$  are two sets of pairwise noncompatible atoms satisfying the condition (F). Let  $\mathcal{M}_1, \mathcal{M}_2$  are families of all atomic blocks of  $E_1$  and  $E_2$  respectively. Then the condition (i) implies the conditions (ii) and (iii):*

- (i)  $P_{F_1}$  and  $P_{F_2}$  are isomorphic generalized effect algebras.
- (ii) There exist bijections  $\psi : \text{At}(E_1) \rightarrow \text{At}(E_2)$  and  $\alpha : \mathcal{M}_1 \rightarrow \mathcal{M}_2$  such that for any  $M \in \mathcal{M}_1$  there exists an isomorphism  $\varphi_M : M \rightarrow \alpha(M)$  of atomic MV-effect algebras such that  $\varphi_M(a) = \psi(a)$  for all  $a \in \text{At}(M)$  and  $\psi(F_1) = F_2$ .
- (iii) There is a bijection  $\psi : \text{At}(E_1) \rightarrow \text{At}(E_2)$  such that
  - (a)  $\psi(F_1) = F_2$ .
  - (b)  $\text{ord}(\psi(a)) = \text{ord}(a)$  for all  $a \in \text{At}(E_1)$ .
  - (c)  $a \longleftrightarrow b$  iff  $\psi(a) \longleftrightarrow \psi(b)$ , for all atoms  $a, b \in \text{At}(E_1)$ .

Moreover, any of the conditions (i) and (ii) implies the condition (iii) and if  $E_1$  and  $E_2$  are complete effect algebras then all the conditions (i), (ii) and (iii) are mutually equivalent.

## REFERENCES

- [1] E.G. Beltrametti, G. Cassinelli, The Logic of Quantum Mechanics, Addison-Wesley, Reading, MA, 1981.

- [2] C.C. Chang, Algebraic analysis of many-valued logics, *Trans. Amer. Math. Soc.* 88 (1958) 467–490.
- [3] D.J. Foulis, M.K. Bennett, Effect algebras and unsharp quantum logics, *Found. Phys.* 24 (1994), 1325–1346.
- [4] J. Hedlíková, S. Pulmannová, Generalized difference posets and orthoalgebras, *Acta Math. Univ. Comenianae* 45 (1996) 247–279.
- [5] J. Paseka, Z. Riečanová, On atoms based isomorphism theorems of lattice or generalized prelattice effect algebras, preprint, 2007.
- [6] Z. Riečanová: Effect algebraic extensions of generalized effect algebras and two-valued states, *Fuzzy Sets and Systems*, to appear.
- [7] Z. Riečanová, The existence of states on every Archimedean atomic lattice effect algebra with at most five blocks, *Kybernetika*, to appear.

DEPARTMENT OF MATHEMATICS AND STATISTICS, FACULTY OF SCIENCE, MASARYK UNIVERSITY,  
JANACKOVO NAM. 2A, 602 00 BRNO, CZECH REPUBLIC  
*E-mail address:* [jpaseka@gmail.com](mailto:jpaseka@gmail.com)